

Selector for Structural Isomers of Neutral Molecules

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(Received 30 December 2007; published 2 April 2008)

We have selected and spatially separated the two conformers of 3-aminophenol (C_6H_7NO) present in a molecular beam. Analogous to the separation of ions based on their mass-to-charge ratios in a quadrupole mass filter, the neutral conformers are separated based on their different mass-to-dipole-moment ratios in an ac electric quadrupole selector. For a given ac frequency, the individual conformers experience different focusing forces, resulting in different transmissions through the selector. These experiments demonstrate that conformer-selected samples of large molecules can be prepared, offering new possibilities for the study of gas-phase biomolecules.

DOI: [10.1103/PhysRevLett.100.133003](https://doi.org/10.1103/PhysRevLett.100.133003)

PACS numbers: 37.20.+j, 33.15.-e, 87.15.-v

During the past decades, the properties of biomolecules in the gas phase have been studied in ever greater detail [1–3]. Although the study of biomolecules outside of their natural environment was met with skepticism in the beginning, spectroscopic studies on isolated species in a molecular beam have proven to be very powerful to understanding their intrinsic properties. Moreover, their native environment can be mimicked by adding solvent molecules one by one. These studies on well-defined biomolecular systems are particularly relevant to benchmark theoretical calculations. Even in the cold environment of a molecular beam, biomolecules exist in various conformational structures. The existence of multiple conformers (structural isomers) has been observed in the study of glycine for the first time [4] and in numerous experiments since then. In many cases, the individual conformers are identified via their different electronic spectra [5,6]. This has been exploited in multiple-resonance techniques to measure, for instance, conformer-specific infrared spectra from which the conformational structures can be deduced [7,8]. Apart from this information on the local minima on the potential energy surface, information on the barriers separating the conformers has been obtained in sophisticated multiple-resonance experiments as well [9].

The preparation of conformer-selected samples of molecules could enable a new class of experiments to be performed on these systems, e.g., electron and x-ray diffraction [10,11] and tomographic imaging experiments [12]. Also, ultrafast dynamics studies on the ground-state potential energy surface would benefit from the availability of these pure samples. For charged species the separation of structurally different molecules has been demonstrated using ion mobility in drift tubes [13,14]. For neutral molecules, no such separation method exists. It has been demonstrated that the abundance of the conformers in the beam can be partly influenced by selective over-the-barrier excitation in the early stage of the expansion [15] or by changing the carrier gas [16]. These methods, however, are not generally applicable nor able to specifically select each of the conformers.

Polar molecules experience a force in an inhomogeneous electric field given by the negative gradient of the Stark energy. If the molecule is in an eigenstate whose Stark energy increases with increasing electric field, a so-called low-field-seeking (LFS) state, it feels a force towards regions of low electric field. Molecules in LFS states can be focused using static inhomogeneous electric fields, and these molecules have been slowed down and trapped using time-varying electric fields [17,18]. Large molecules have a high density of rotational states, and due to the interaction between these states all of them are high-field seeking (HFS), i.e., they feel a force towards regions of high electric field. Molecules in HFS states can be dynamically focused using the alternating gradient principle [19]. This has been demonstrated for ammonia in HFS states [20,21]. It has also been applied in the deceleration of CO [22,23], YbF [24], and benzonitrile [25] molecules in HFS states, as well as in the ac trapping of ND_3 [26].

The conformers of a specific biomolecule all have the same mass m and the same connectivities between the atoms (constitution) but differ by the orientations of their functional groups in the molecular frame, i.e., by their folding pattern. The vectorial sum of the local dipole moments of the functional groups largely determines the overall dipole moment of the molecule. The different dipole moments μ of the conformers can be exploited to select individual conformers using dynamic focusing with ac electric fields. This is most easily implemented in a setup using high-voltage electrodes in a quadrupole arrangement around a molecular beam. The operation principle of such an m/μ selector is equivalent to that of the m/q quadrupole mass filter for charged particles, where q is the charge [27]. In the m/μ selector the most polar quantum states of a given conformer are focused most efficiently and have the highest transmission. These are also the quantum states that can be aligned or oriented best, using intense laser fields or strong static electric fields, respectively. Therefore, experiments that rely on highly oriented samples, such as tomographic or diffraction imaging experiments, would particularly benefit from the conformer-selected polar ensembles generated by the m/μ selector.

In this Letter, we demonstrate the selective transmission of the cis and trans conformers of neutral 3-aminophenol (C_6H_7NO) through an ac quadrupole m/μ selector. The cis and trans conformers of 3-aminophenol are used here as prototypes for the different structural isomers of biomolecules. The rotational envelopes of the electronic origin transitions of the individual conformers provide information on the rotational state specific transmission of the device.

The experimental setup is shown in Fig. 1(a). A sample of 3-aminophenol is heated to 110°C and coexpanded in 2 bar of Kr through a pulsed nozzle operated at a repetition rate of 10 Hz. The mean velocity of the molecules in the beam is approximately 480 m/s with a velocity spread (full width at half maximum) of about 10%. After passing two skimmers, placed 5 and 15 cm downstream from the nozzle, the molecules enter a second, differentially pumped vacuum chamber, in which the m/μ selector is placed. The selector consists of four polished, 1 m long cylindrical stainless steel electrodes of 4 mm diameter. High voltages of 12 kV against ground are applied as shown in Fig. 1(b). The gaps are 0.9 mm between two adjacent electrodes and 3 mm between two opposing electrodes, resulting in a field strength of 45 kV/cm on the center line and a maximum field strength of 135 kV/cm. Using three high-voltage switches, the field is rapidly switched ($<1 \mu\text{s}$) between the two electric field configurations shown in Fig. 1(b). Switching from one configuration to the other interchanges the role of the x and y axes; i.e., it interchanges the directions of transverse focusing and defocusing as indicated by the white arrows in Fig. 1(b). The resulting dynamic focusing of neutral mole-

cules, which is very similar to the dynamic focusing of charged particles [29], has been described in detail elsewhere [19,23].

The transmitted 3-aminophenol molecules are ionized 1.21 m downstream from the nozzle using two-color resonance-enhanced multiphoton ionization, $(1+1')$ -REMPI, as shown in Fig. 1(c). Prior to entering the detection region, the molecules have to pass through a 2 mm diameter aperture positioned on the center line of the selector. The axis of the m/μ selector is tilted against the axis of the incoming molecular beam by 0.3° , such that predominantly molecules that are transported through the selector enter the detection region. The laser beams used for excitation and ionization are unfocused, and have diameters of 4 and 2 mm, respectively. To minimize saturation effects, the energy of the frequency-doubled pulsed dye laser for electronic excitation (290 nm) is reduced to $20 \mu\text{J}/\text{pulse}$, while the ionization laser (355 nm) is operated with 5 mJ/pulse. The ions are mass-selectively detected in a time-of-flight mass spectrometer. The cis and trans conformers can be selectively detected due to their distinct $S_1 \leftarrow S_0$ excitation wave numbers of $34\,108 \text{ cm}^{-1}$ and $34\,467 \text{ cm}^{-1}$, respectively.

From the precisely known rotational constants and dipole moments [30] the energies of the rotational states of cis-3-aminophenol and trans-3-aminophenol are calculated as a function of electric field strength. Figure 2 shows the resulting Stark curves for the lowest rotational states of both species. The transmission characteristics of the selector depend on the effective dipole moment μ_{eff} (the negative of the slope of the Stark curve), the electric field gradients, and the ac frequency. Similar to the frequency dependence in quadrupole mass filters, molecules with a given value of μ_{eff} are only transmitted through the selector within a finite range of frequencies. At too low frequencies molecules are deflected and lost in one transverse dimension before they are refocused. For high frequencies the time-averaged potential becomes flat resulting in a strongly reduced transmission. The ac frequency for optimum transmission increases with increasing μ_{eff} . When a constant ac frequency is applied, a μ_{eff} selection is performed.

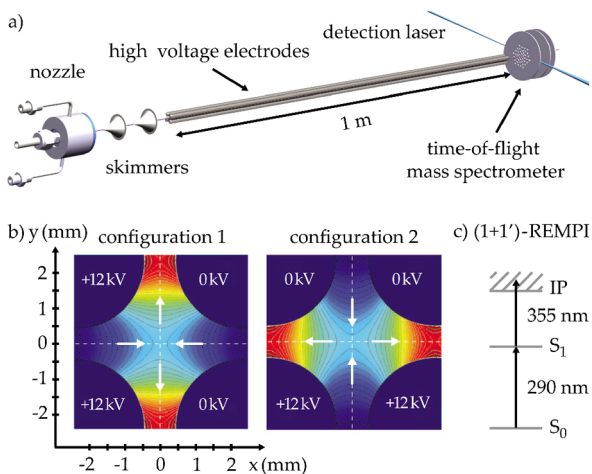


FIG. 1 (color online). (a) Scheme of the experimental setup. (b) Cut through the high-voltage electrodes and electric field configurations; gray (red) areas represent regions of strong electric field, black (blue) areas regions of weak electric field. The electric field strength at the center is 45 kV/cm and contour lines are given every 4 kV/cm. White arrows indicate the direction of the force experienced by molecules in high-field seeking quantum states. (c) Two-color ionization detection scheme.

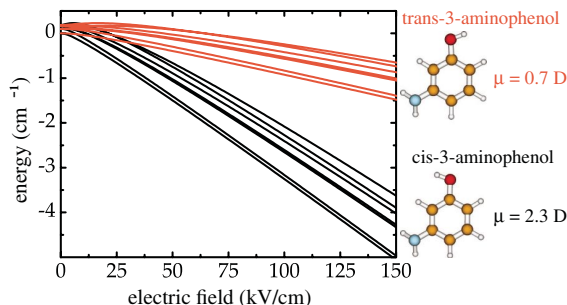


FIG. 2 (color online). Molecular structures, dipole moments, and energies of the lowest rotational states of cis- and trans-3-aminophenol as a function of the electric field strength.

Scans of the ac frequency for *cis*-3-aminophenol and *trans*-3-aminophenol are shown in Fig. 3. The transmission measurements are performed with the excitation laser frequency set close to the band origin of the respective conformer. The ac frequency is scanned from 0 to 5 kHz in steps of 50 Hz. For a given conformer, the number of transmitted molecules is measured with and without applied high voltages, and the ratio of these two measurements is plotted in Fig. 3. The start phase of the switching cycle determines the overall transmission of the selector. In all measurements we start with a half-period of focusing along the horizontal axis (configuration 1). As a consequence, the phase of the switching cycle at the moment that the molecules exit the selector changes with the applied frequency. For *cis*-3-aminophenol a clear enhancement of the transmission is observed for ac frequencies in the range from 2–3.5 kHz, whereas for *trans*-3-aminophenol a weaker transmission maximum is observed around 1.5 kHz. Both the higher frequency and the higher transmission for *cis*-3-aminophenol reflect its considerably larger dipole moment compared to that of the *trans* conformer. The central dip in the transmission curve at 2.7 kHz is due to effects of the exact phase of the ac switching cycle at the exit of the selector. This phase determines the shape of the molecular packet in the detection region, and thereby its overlap with the laser beams.

As discussed above, the transmission of the selector depends on the effective dipole moment μ_{eff} of the individual quantum states. Although we cannot detect individual rotational states, there is a well-defined relation between the excitation laser frequency and the rotational states that are probed. In Fig. 4 the $(1 + 1')$ -REMPI spectrum of 3-aminophenol is shown. In the center of each of the vibronic bands, predominantly transitions from low- J states are probed, whereas the wings of the rotational en-

velopes contain mostly transitions from high- J states. The inset of Fig. 4 shows the rotational contour of the origin transition of *cis*-3-aminophenol on an enlarged wave number scale measured with and without electric fields, for different ac frequencies. For a frequency of 1.6 kHz the wings of the rotational envelope are increased, whereas the intensity of the central part of the band is actually decreased. This directly reflects that for this ac frequency high- J states are efficiently transported through the selector, whereas low- J states, that generally have a larger μ_{eff} , are over-focused and have a lower transmission. For a frequency of 2.95 kHz the whole rotational envelope is clearly increased. The largest enhancement is now observed for the central part of the rotational envelope, where mostly low- J states are probed. For *trans*-3-aminophenol similar, albeit weaker, changes in the rotational envelope are observed.

Monte Carlo trajectory calculations are performed to simulate the transmission curves for fixed excitation laser frequencies. For this, the rotationally resolved electronic excitation spectrum is calculated using the known rotational constants and transition moment orientations [30,31]. A rotational temperature of 4 K yields a rotational envelope that agrees best with the observations. A rectangular spectral profile of the laser with a width of 0.15 cm^{-1} is assumed. For all rotational states that are probed within this bandwidth of the laser, Monte Carlo simulations are performed, and individual transmission curves are calculated. From the calculated line strengths and populations, a weight for every single quantum state is determined. The weighted sum over the individual transmission curves is shown together with the experimental data in Fig. 3 (dashed lines). These simulations nicely reproduce the peak position and the low-frequency cutoff of the experimental transmission curves. On the high-frequency side the experimentally observed transmission decreases faster

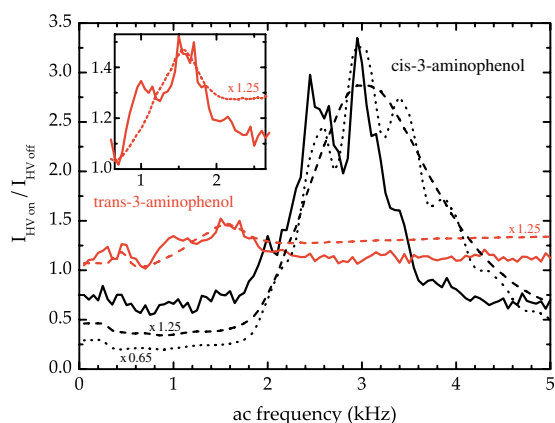


FIG. 3 (color online). Experimental (solid lines) and simulated (dashed lines) transmission as a function of ac frequency for *cis*-3-aminophenol (black) and *trans*-3-aminophenol [gray (red)]. The inset shows an enlarged view of the transmission curve of *trans*-3-aminophenol. The black dotted line represents simulations for *cis*-3-aminophenol that include the details of the detection process.

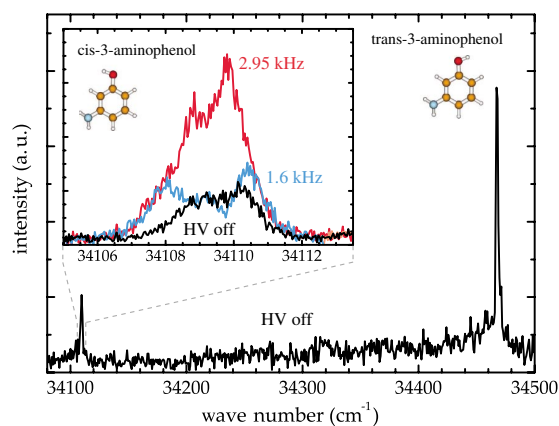


FIG. 4 (color online). $(1 + 1')$ -REMPI spectrum of 3-aminophenol. The spectrum is measured without high voltage (HV) applied to the selector, and shows the two origin transitions. In the inset, the rotational envelope of the *cis*-3-aminophenol origin transition is shown without high voltage (black) and with high voltage for ac frequencies of 2.95 kHz [dark gray (red)] and 1.6 kHz [light gray (blue)].

than predicted, which we attribute to mechanical misalignment. Taking into account the phase-dependent shape of the molecular packet and its spatial overlap with the detection laser beams, the observed modulation of the transmission peak for cis-3-aminophenol is correctly reproduced (dotted line).

In the experiments presented here, the selector is operated under conditions for optimum transmission, equivalent to the “rf-only” operation mode of quadrupole mass filters. The resolution $\mu_{\text{eff}}/\Delta\mu_{\text{eff}}$ of the selector is only about two in this case. In m/q filters, the resolution $m/\Delta m$ is increased by adding a dc offset to the rf potentials, at the cost of a reduced transmission. In the m/μ selector a better resolution can be achieved by adding a static defocusing field to the two configurations of the electric field that we use here. This can be achieved by using different high voltages for the two electric field configurations or, more easily, by changing the duty cycle, i.e., by applying the presently used field configurations for different time intervals. For many biomolecules, e.g., amino acids and peptides, the various conformers have large and widely different dipole moments. For phenylalanine, for instance, at least six conformers have been observed [7,8] and their dipole moments are calculated to range from 1 to 5.5 D. Therefore, the selection of its conformers would be feasible even at the present resolution.

In summary, we have selected and separated the cis and trans conformers of 3-aminophenol using switched electric fields in a quadrupole m/μ selector. The conformers are separated based on their distinct frequency dependent transmission characteristics. The dynamic focusing works best for the states with the largest effective dipole moments, which are the lowest rotational states. Here, molecular packets with an excess of cis-3-aminophenol in low rotational states are created at an ac frequency of 2.95 kHz. Such conformer-selected molecular packets offer interesting perspectives for a variety of experiments. Since the selected states are the most polar ones, these samples are particularly useful for experiments in which aligned or oriented molecules are desired.

Financial support from the Deutsche Forschungsgemeinschaft within the priority program 1116 “Interactions in ultracold gases” is acknowledged.

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